



DØnote 4524-CONF

## Reconstruction of Semileptonic Lambda-b Decays

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URL <http://www-d0.fnal.gov>  
(Dated: August 26, 2004)

The decay  $\Lambda_b^0 \rightarrow \Lambda_c^+ \mu^- \bar{\nu}_\mu X$  is reconstructed in data corresponding to approximately  $350 \text{ pb}^{-1}$  collected by the DØ detector between April 2002 and May 2004.

*Preliminary Results for Summer 2004 Conferences*

## I. INTRODUCTION

This analysis has so far involved searching for semileptonic decays of  $\Lambda_b$  baryons in data taken by the DØ detector in Run II. These decays can be used along with decays of  $B_d$  mesons to measure the ratio of lifetimes of these two hadrons. The decays are due to the decay of the  $b$  quark, but the particular environment of light quarks and gluons also has an effect on the lifetime. The differences in lifetimes of hadrons containing a heavy quark are predicted to vanish as the mass  $m_Q$  of the heavy quark tends to infinity [1]. These effects have been predicted for  $B$  hadrons using a technique known as Operator Product Expansion (OPE). Using OPE the lifetime ratios of  $B$  hadrons have been predicted and are in reasonable agreement with experimentally measured values, except in the case of the  $\Lambda_b$ . For the ratio  $\tau(\Lambda_b)/\tau(B_d^0)$  theoretical predictions, given their uncertainties, strongly favor values above 0.9, whereas the experimental evidence suggests values close to 0.8, with errors of approximately 0.05. Data taken at DØ at the high luminosity provided by the Fermilab Tevatron in Run II should allow a large  $\Lambda_b$  sample to be obtained and therefore allow a high precision measurement of this ratio, using a method similar to that described in [2].

The  $\Lambda_b$  decay searched for in this analysis is  $\Lambda_b^0 \rightarrow \Lambda_c^+ \mu^- \bar{\nu}_\mu X$ , where  $X$  is any other particle. Charge conjugates are implied throughout this paper. This decay has the advantage of a relatively high branching fraction, since the branching fraction for  $\Lambda_b^0 \rightarrow \Lambda_c^+ l^- \bar{\nu}_l X$  is  $7.7 \pm 1.8\%$  [3]. The selections were designed to reconstruct the  $\Lambda_c$  in the decays  $\Lambda_c^+ \rightarrow K_s^0 p^+$  and  $\Lambda_c^+ \rightarrow \Lambda \pi^+$ . Decays were reconstructed using data from approximately  $350 \text{ pb}^{-1}$  integrated luminosity, acquired by the DØ detector between April 2002 and May 2004. This note outlines the selections used to isolate the  $\Lambda_b$  decays and shows the signals obtained.

## II. EVENT SELECTION

For each event the primary vertex was determined using the method described in [4] and the charged particles were clustered into jets using the DURHAM clustering algorithm [5]. Muons were selected using the standard DØ tools and were then also required to have at least two track segments in the muon chambers associated to a central track and a value of  $p_T$  of at least  $2.0 \text{ GeV}/c$ . The products from the decay of  $\Lambda_c$  were then searched for among tracks belonging to the same jet as an identified muon.

### A. Selection of $\Lambda$ and $K_s$

The  $\Lambda$  and  $K_s$  were reconstructed as  $V^0$  decays to two oppositely charged tracks. These tracks were required to form a secondary vertex with  $\chi^2$  per degree of freedom of less than 25 and a distance  $d_{V^0}$  from the primary vertex of at least  $4\sigma(d_{V^0})$  if  $\sigma(d_{V^0})$  was less than  $0.5 \text{ cm}$ . The two tracks were required to have no more than two hits in the tracking system between them before the secondary vertex and each track was allowed no more than one miss after the reconstructed vertex. A loose cut on the direction of the  $V^0$  was also applied. The angle  $\alpha_{V^0}$  between the momentum of the reconstructed particle and the vector from the primary to the secondary vertex in the axial plane was required to fulfill the condition  $\cos(\alpha_{V^0}) > 0.8$ . The axial  $\epsilon_T$  and stereo  $\epsilon_L$  impact parameter projections of each track with respect to the primary vertex and their errors were calculated and the combined significance  $(\epsilon_T/\sigma(\epsilon_T))^2 + (\epsilon_L/\sigma(\epsilon_L))^2$  was required to be greater than 9 for both tracks and greater than 16 for at least one track. Further selections were then applied to the tracks and reconstructed vertex. Any vertices containing the identified muon as one of the tracks were rejected and both tracks were required to have at least one hit in the Central Fiber Tracker. The particle reconstructed was required to have  $p_T$  of at least  $0.7 \text{ GeV}/c$ .

To search for decays of  $\Lambda$  the values for the mass of the  $V^0$  when the decay products were assumed to be both permutations of proton and pion were calculated. These are plotted in Figure 1 for a small subsample of the events, where the peak due to  $\Lambda$  can clearly be seen. Similar plots were used to set a mass range of  $1.109 \text{ GeV}/c^2 - 1.120 \text{ GeV}/c^2$  for accepted  $\Lambda$  candidates. The width of the peak is approximately  $2.8 \text{ MeV}$  so these limits correspond to approximately  $1.75\sigma$  below the central mass of  $1.115 \text{ GeV}/c^2$  and  $2.14\sigma$  above. If the mass of the combination with the proton having higher momentum was within this range, the  $V^0$  was accepted as a  $\Lambda$  and assigned this mass. For candidates not accepted the mass of the combination with a higher momentum pion was compared with the mass range, and accepted if it was within this interval. The charge of the proton was used to determine whether the reconstructed candidate was a  $\Lambda$  or  $\bar{\Lambda}$ .

To select  $K_s$  the decay products of the  $V^0$  were assumed to be  $\pi^+\pi^-$  and its invariant mass was calculated. A typical plot of this mass distribution for a small sample of events is shown in Figure 2. Here the peak due to  $K_s$  is also clearly seen, at a central mass of  $0.494 \text{ GeV}/c^2$ . Using similar plots the mass interval that would be accepted as

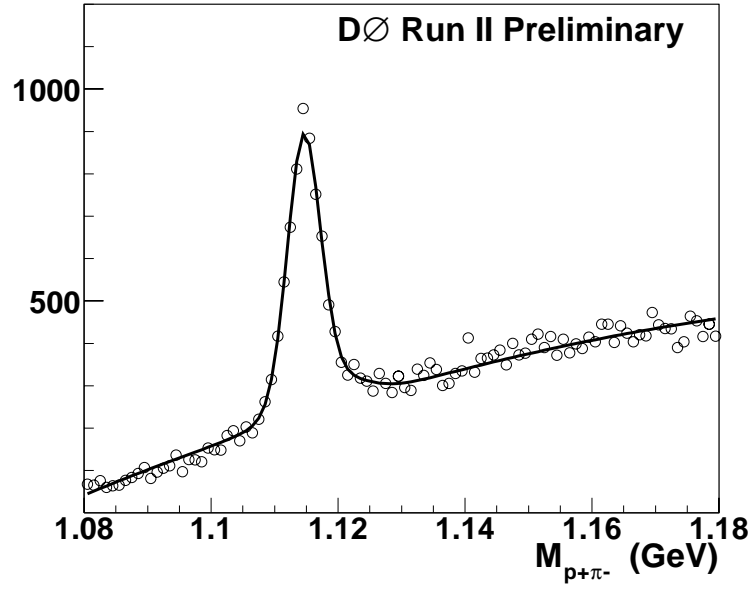


FIG. 1: The invariant mass of  $V^0$  for decays to  $p\pi$  in a small subsample of the data, showing the peak due to  $\Lambda$  at a mass of  $1115 \text{ MeV}/c^2$ .

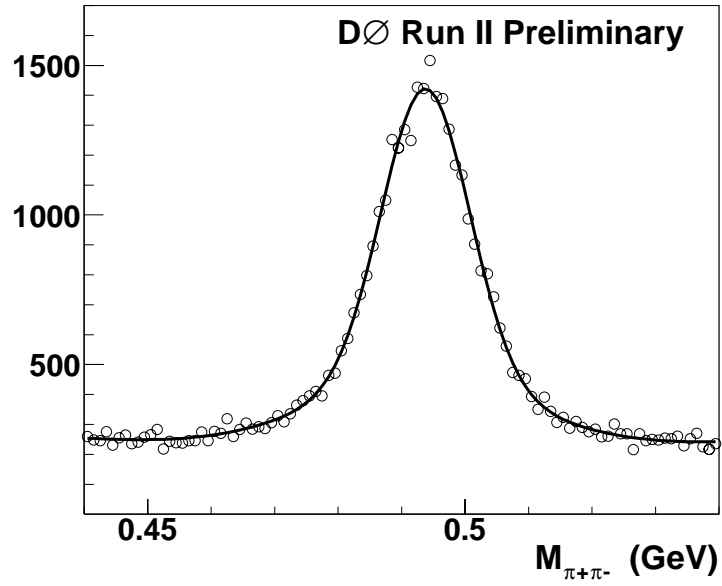


FIG. 2: The invariant mass of  $V^0$  for decays to  $\pi^+\pi^-$  in a small subsample of the data, showing the peak due to  $K_s$  at a mass of  $494 \text{ MeV}/c^2$ .

$K_s$  was set at  $0.480 \text{ GeV}/c^2 - 0.5075 \text{ GeV}/c^2$ . The width of this peak is approximately  $7.6 \text{ MeV}$  so these limits are both approximately  $1.8\sigma$  away from the central mass.

If the mass conditions for  $\Lambda$  and  $K_s$  were both satisfied the  $V^0$  was rejected.

TABLE I: Final cuts used for the two decay channels

	$\Lambda_c \rightarrow \Lambda \pi$	$\Lambda_c \rightarrow K_s p$
$\cos(\alpha_T^c)$	$> 0.84$	$> 0.91$
$d_T^c/\sigma(d_T^c)$	$> 3.6$	$> 3.8$
$M(\Lambda_c/D + \mu)$	$3.1\text{--}5.5 \text{ GeV}/c^2$	$3.4\text{--}5.4 \text{ GeV}/c^2$
$P_T^{\Lambda_c}(\mu)$	$> 0.35 \text{ GeV}/c$	$> 0.35 \text{ GeV}/c$
$d_T^{bc}/\sigma(d_T^{bc})$	$< 3.2$	$< 3.3$
$(\epsilon_T/\sigma(\epsilon_T))^2 + (\epsilon_L/\sigma(\epsilon_L))^2$	$> 3.2$	$> 4.6$

### B. Reconstruction of $\Lambda_c$

For each reconstructed  $\Lambda$  or  $K_s$ , the remaining tracks in the jet were searched for the proton or pion produced in the decay of  $\Lambda_c$ . Each track with  $P_T > 1.0 \text{ GeV}/c^2$  and at least 2 hits in the silicon detector was formed into a vertex with the  $\Lambda$  or  $K_s$ . The reconstructed vertex was required to have a value of  $\chi^2$  per degree of freedom less than 9. The reconstructed  $\Lambda_c$  candidate was then required to form a vertex with the muon. This vertex was also required to have  $\chi^2$  per degree of freedom less than 9. Reconstructed decay chains in which the  $\Lambda_c$  decay vertex was closer to the primary vertex than the  $\Lambda_b$  vertex were also rejected. To search for either of the decays it was then only necessary to assign the appropriate particle type to the track that was vertexed with the  $\Lambda$  or  $K_s$  candidate and plot the invariant mass of the new particle created in the vertex.

### C. Selection Criteria

Other variables were calculated and their values used as selection criteria for all channels. The distance  $d_T^c$  in the axial plane between the  $\Lambda_c$  decay vertex and the primary vertex and the error on this distance  $\sigma(d_T^c)$  were calculated and  $d_T^c/\sigma(d_T^c)$  was used as a discriminating variable. Since the  $\Lambda_c$  candidate should be produced in a decay of a  $\Lambda_b$  and the  $b$  hadrons have relatively long lifetimes compared to charmed hadrons, this distance should be significant. The angle  $\alpha_T^c$  in the axial plane between the flight direction of the  $\Lambda_c$  candidate and the vector from the primary vertex to the  $\Lambda_c$  decay vertex was also calculated and a minimum value was set for the  $\cos(\alpha_T^c)$ . The transverse momentum  $P_T^{\Lambda_c}(\mu)$  of the muon with respect to the direction of the  $\Lambda_c$  candidate was computed and cuts requiring a minimum value of this variable were used. The muon was also combined with the  $\Lambda_c$  candidate into a new particle, and the invariant mass of this combination was calculated and upper and lower limits were set for this mass. Due to the short lifetime of the  $\Lambda_c$ , a limit on the distance  $d_T^{bc}$  between the  $\Lambda_b$  and  $\Lambda_c$  decay vertices divided by its error was imposed. The combined impact parameter significance  $(\epsilon_T/\sigma(\epsilon_T))^2 + (\epsilon_L/\sigma(\epsilon_L))^2$  of the  $\Lambda_c$  was calculated and events with this parameter below a certain value were rejected.

The cuts were optimized separately for each channel, the values used are shown in Table I. Some of these cuts will bias the lifetime distributions but in the case of measurement of lifetime ratios [2] this bias should not affect the final measurement.

## III. RESULTS

### A. $K_s$ Events

For the events in which a  $K_s$  was identified the mass for the  $K_s p$  hypothesis is plotted in Figure 3. The upper points show the reconstructed events in which the muon and  $D$  candidate have opposite charges as is expected in the decay  $\Lambda_b^0 \rightarrow \Lambda_c^+ \mu^- \bar{\nu}_\mu X$ . The lower histogram shows the wrong sign combination for the muon and the  $\Lambda_c$  candidate. Here a peak close to the PDG [3] value of the  $\Lambda_c$  mass of  $2284.9 \pm 0.6 \text{ MeV}/c^2$  is clearly observed in the plot for right sign events. This peak is almost  $9\sigma$  above the background so this is very strong evidence that a signal is observed. However, semileptonic decays of  $B$  mesons will also be selected, and the reflection of these events in the  $M(K_s p)$  mass spectrum may contribute to this peak. It is expected that the plots also contain contributions from the decays  $\bar{B}_d^0 \rightarrow D^+ \mu^- \bar{\nu}_\mu X$  with  $D^+ \rightarrow K_s^0 \pi^+$  and  $\bar{B}_s^0 \rightarrow D_s^+ \mu^- \bar{\nu}_\mu X$  with  $D_s^+ \rightarrow K_s^0 K^+$ . A peak due to  $D$  and  $D_s$  from these decays is observed in the region between  $1.81 \text{ GeV}/c^2$  and  $1.92 \text{ GeV}/c^2$  if the mass of the  $K_s \pi$  hypothesis is plotted. To check that reflection from this peak is not the cause of the peak in Figure 3, the events with  $M(K_s \pi)$  between

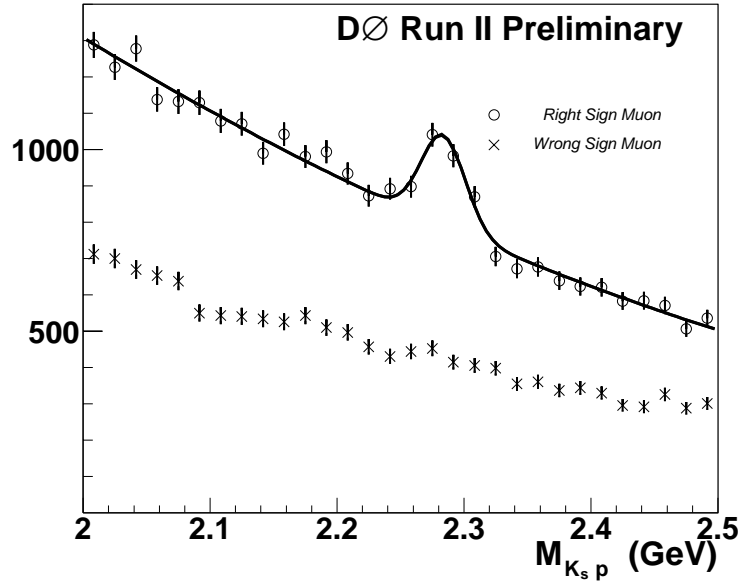


FIG. 3: The invariant mass of  $K_s p$  for reconstructed decays containing a  $K_s$ , for right and wrong charge correlation with the muon. The fitted peak is at a mass of  $2284 \pm 2.2(stat)$  MeV/ $c^2$ .

1.81 GeV/ $c^2$  and 1.92 GeV/ $c^2$  which survived the cuts were selected and the mass of the  $K_s p$  hypothesis was plotted. This is shown in Figure 4, where it can be seen that these events are distributed over a wide mass range and the main peak is at a lower mass than the  $\Lambda_c$ , so the  $B$  meson decays mainly contribute to the background observed in Figure 3. A small peak of around 70 events is observed around 2.28 GeV/ $c^2$  in Figure 4. This may be due to reflection of  $D$  and  $D_s$  events in the  $\Lambda_c$  mass range. However, this peak contains a small fraction of the events observed in the  $\Lambda_c$  signal, so the majority of the events in the peak in Figure 3 are due to  $\Lambda_b$  decays.

### B. $\Lambda$ Events

For events containing a  $\Lambda$  the mass  $M(\Lambda\pi)$  is shown in Figure 5. Besides the peak close to 2.28 GeV/ $c^2$  as expected, there is also an indication of another peak at a slightly lower mass. Therefore this plot has been fitted with the sum of two Gaussians and a quadratic background, where the widths of the two Gaussians were allowed to vary independently of each other. The peak at higher mass is approximately  $6.3\sigma$  above the background and the lower peak is approximately  $2.9\sigma$ . There is strong evidence for the signal of  $\Lambda_c^+ \rightarrow \Lambda\pi^+$ , but significantly higher statistics would be needed to confirm the existence of a signal for the lower peak. The lower peak is thought to be due to the decay  $\Lambda_c^+ \rightarrow \Sigma^0\pi^+$ .  $\Sigma^0$  decays to  $\Lambda\gamma$  close to 100% of the time, and has a very short lifetime, so it is likely that the  $\Lambda$  and pion have been found and reconstructed into a vertex for many of these decays. Since the photon is not included, these decays should cause a peak to be seen at about 74.5 MeV/ $c^2$  below the mass of the  $\Lambda_c$  [3]. The mass difference shown on the plot is  $78 \pm 11.4(stat)$  MeV/ $c^2$ . The branching fraction of this decay is given by the PDG [3] to be  $(9.9 \pm 3.2) \cdot 10^{-3}$  whereas the decay to  $\Lambda\pi$  has branching fraction  $(9.0 \pm 2.8) \cdot 10^{-3}$ , but using these selections a lower number of decays to  $\Sigma^0\pi$  are observed. It should be also possible to use these decays in the later analysis, to increase the statistics for the  $\Lambda_b$  sample.

## IV. CONCLUSIONS

Using the selections outlined in earlier sections and approximately 350 pb $^{-1}$  of data taken by DØ in Run II the signal for semileptonic decays of  $\Lambda_b$  has successfully been isolated using reconstruction of  $\Lambda_c$  in the modes  $\Lambda_c^+ \rightarrow K_s^0 p^+$  and  $\Lambda_c^+ \rightarrow \Lambda\pi^+$ . An indication of the decay of  $\Lambda_c^+ \rightarrow \Sigma^0\pi^+$  has also been observed. This sample will be used in a

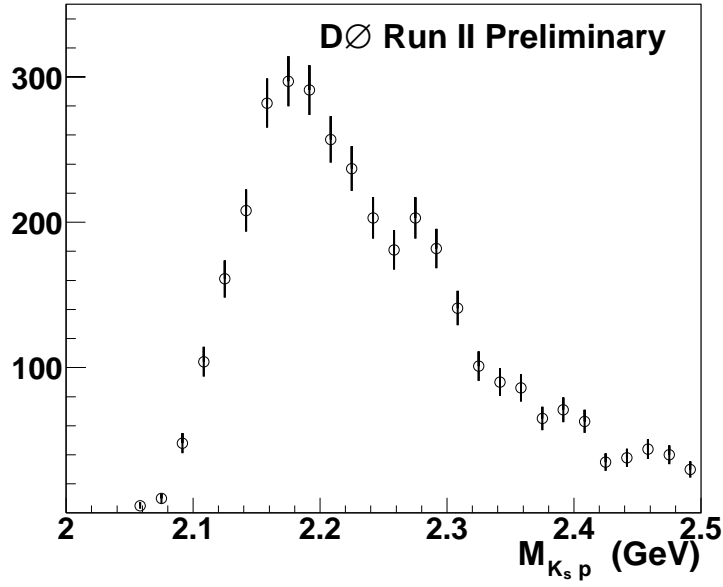


FIG. 4: Invariant mass of  $K_s^0 p$  for right sign events passing the selection cuts used for  $\Lambda_c \rightarrow K_s p$  with  $M(K_s \pi)$  between  $1.81 \text{ GeV}/c^2$  and  $1.92 \text{ GeV}/c^2$ .

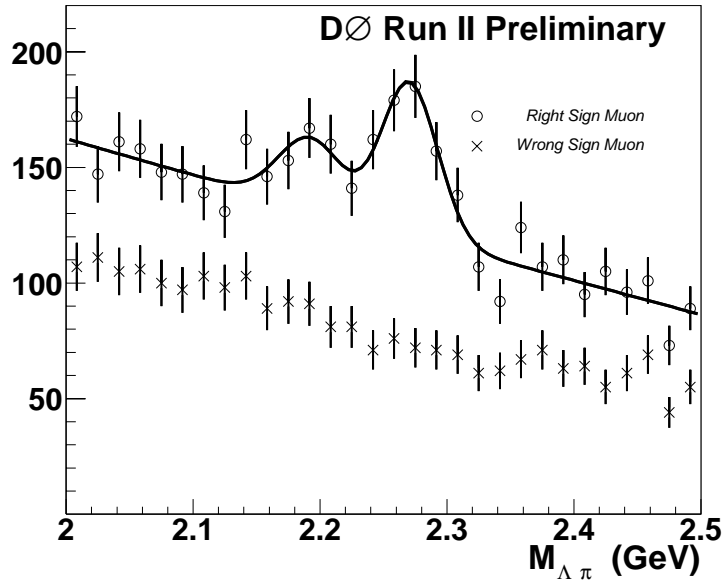


FIG. 5: Invariant mass of  $\Lambda \pi$  for reconstructed decays containing a  $\Lambda$ , for both right and wrong sign charge correlation with the muon. The peak at higher mass is at  $2271 \pm 5.0(stat) \text{ MeV}/c^2$ , the lower mass peak is at  $2193 \pm 10.3(stat) \text{ MeV}/c^2$  and is due to the decay  $\Lambda_c^+ \rightarrow \Sigma^0 \pi^+$  with  $\Sigma^0 \rightarrow \Lambda \gamma$ .

measurement of the lifetime ratio of  $\Lambda_b$  and  $B_d$ .

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